

## Comparative Analysis of the Kinetic Characteristics of L-Type Calcium Channels in Cardiac Cells of Hibernators

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**ABSTRACT** An undefined property of L-type  $\text{Ca}^{2+}$  channels is believed to underlie the unique phenotype of hibernating hearts. Therefore, L-type  $\text{Ca}^{2+}$  channels in single cardiomyocytes isolated from hibernating versus awake ground-squirrels (*Citellus undulatus*) were compared using the perforated mode of the patch-clamp technique, and interpreted by way of a kinetic model of  $\text{Ca}^{2+}$  channel behavior based upon the concept of independence of the activation and inactivation processes. We find that, in hibernating ground-squirrels, the cardiac L-type  $\text{Ca}^{2+}$  current is lower in magnitude when compared to awake animals. Both in the awake or hibernating states, kinetics of L-type  $\text{Ca}^{2+}$  channels could be described by a  $d^2f_1f_2$  model with an activation and two inactivation processes. The activation (or  $d$ ) process relates to the movement of the gating charge. The slow (or  $f_1$ ) inactivation is associated with movement of gating charge and is current-dependent. The rapid (or  $f_2$ ) inactivation is a complex process which cannot be represented as a single-step conformational transition induced by the gating charge movement, and is regulated by  $\beta$ -adrenoceptor stimulation. When compared to awake animals, the kinetic properties of  $\text{Ca}^{2+}$  channels from hibernating ground-squirrels differed in the following parameters: (1) pronounced shift (15–20 mV) toward depolarization in the normalized conductance of both inactivation components, and moderate shift in the activation component; (2) 1.5–2-fold greater time constants; and (3) two-fold greater activation gating charge. Thus, L-type  $\text{Ca}^{2+}$  channels apparently switch their phenotype during the hibernating transition. Stimulation of  $\beta$ -adrenoceptors by isoproterenol, reversed the hibernating kinetic- (but not amplitude-) phenotype toward the awake type. Therefore, an aberrance in the  $\beta$ -adrenergic system can not fully explain the observed changes in the L-type  $\text{Ca}^{2+}$  current. This suggests that during hibernation additional mechanisms may reduce the single  $\text{Ca}^{2+}$  channel-conductance and/or keep a fraction of the cardiac L-type  $\text{Ca}^{2+}$  channel population in a non-active state.

## INTRODUCTION

Hibernating animals, such as ground squirrels, hedgehogs or chipmunks, fall periodically into hibernation, which is characterized by a drop in body temperature from 37 to 5°C (Wang, 1978; Snapp and Heller, 1981). However, the drop in body temperature is a secondary phenomenon that follows changes in the primary integrative mechanism(s) of hibernating adaptation (Kondo, 1987; Heldmaier and Ruf, 1992; Kondo and Kondo, 1993). A prominent feature of the hibernating state is that the pronounced reduction in heart rate from 400 to less than 10 beats *per min* does not lead to death contrary to nonhibernating species (Lyman et al., 1982; Duker et al., 1983; Wang, 1978). Thus, hearts of hibernating species possess a unique cardioprotective phenotype with each hibernating cycle.

Whereas in nonhibernating species (e.g., guinea pig) cooling prolongs the plateau phase of the action potential, in hibernating animals (e.g., hedgehog) hypothermia shortens the plateau phase of the cardiac action potential (Duker et al., 1986; Burlington et al., 1989; Smith et al., 1989; Herve et al., 1992). This suggests that cardiac  $\text{Ca}^{2+}$  channels, which participate in the plateau phase of an action

potential, have unique properties in hibernators. Indeed, in papillary muscles from hibernating chipmunks,  $\text{Ca}^{2+}$  current appears to be virtually absent, although in nonhibernating chipmunks this current does not differ from that of nonhibernating species (Kondo, 1986a). Because  $\text{Ca}^{2+}$  influx is essential for excitation-contraction coupling in cardiac muscle (Fabiato and Fabiato, 1979; Fabiato and Baumgarten, 1984), it has been postulated, on the basis of results obtained in multicellular preparations, that the initiation of myocardial contraction in hibernation may be triggered by  $\text{Ca}^{2+}$  influx through the Na/Ca exchanger or controlled through intracellular  $\text{Ca}^{2+}$  release mechanisms (Kondo and Shibata, 1984; Kondo, 1986b; Liu et al., 1991). However, in single cardiac cells during hibernation, we have demonstrated that  $\text{Ca}^{2+}$  current does flow through L-type  $\text{Ca}^{2+}$  channels, albeit dramatically reduced in magnitude, and could participate in the excitation-contraction coupling process during hibernation (Alekseev et al., 1994).

The dramatic reduction in inward  $\text{Ca}^{2+}$  current in hibernation may be due to several mechanisms including a decrease in the number of functional voltage-dependent  $\text{Ca}^{2+}$  channels (Kondo, 1986a), and/or a decrease in the cAMP-dependent phosphorylation of  $\text{Ca}^{2+}$  channels resulting from the decrease in the levels of circulating catecholamines (Gainullin and Saxon, 1988; Saxon and Gainullin, 1989). Yet, no definite characterization of the biophysical properties and regulation of voltage-dependent  $\text{Ca}^{2+}$  channels during hibernation has been performed. This is due in part

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to the fact that studies have been conducted thus far almost exclusively in multicellular preparations or under hypothermic (but not under hibernating) conditions. Comparative analysis of the electrophysiological characteristics of  $\text{Ca}^{2+}$  channels during the hibernating versus the nonhibernating state are lacking. Herein, we analyzed the kinetic properties of L-type  $\text{Ca}^{2+}$  channels in cardiomyocytes isolated from ground squirrels during the hibernating and nonhibernating states, and establish the role of  $\beta$ -adrenoceptor-mediated regulation of L-type  $\text{Ca}^{2+}$  channels during hibernation.

## MATERIALS AND METHODS

### Isolation of single cardiac cells

Ground squirrels (*Citellus undulatus*) trapped at the end of summer in the Yakutia region of Siberia were kept under conditions close to natural. Experiments were performed on animals in two different physiological states: hibernation or spontaneous arousal (active). The body temperature of animals was 6–7.5°C during hibernation, and 37°C during spontaneous arousal. Single ventricular cardiocytes were isolated by the pronase/collagenase method described previously (see Alekshev et al., 1994).

### Current recording

Currents were recorded by the perforated patch-clamp technique (Korn et al., 1991). Isolated cardiocytes were placed in a recording chamber and bathed with the following (in mM): 20 glucose, 5  $\text{MgCl}_2$ , 2  $\text{CaCl}_2$ , 10 HEPES, 80 NaCl, 20 tetraethylammonium chloride (TEA-Cl), 1.2  $\text{KH}_2\text{PO}_4$ , and 10 CsCl (pH 7.3). Pipettes fabricated from soft molybdenum glass had a resistance of 1–5 M $\Omega$  and were filled with the following (in mM): 130 CsCl, 5  $\text{MgCl}_2$ , and 10 HEPES (pH 7.3). The outward potassium current was blocked by TEA, and by replacing  $\text{K}^+$  with  $\text{Cs}^+$ . Solutions were passed through a filter with a diameter of 0.22 mm (Millipore, Boston, MA). Amphotericin B (0.5 mg; Sigma Chemical Co., St. Louis, MO), which served as a perforating compound, was dissolved in 9  $\mu\text{l}$  of dimethylsulfoxide and added to the pipette solution. The final concentration of the antibiotic was 200–250  $\mu\text{g}/\text{ml}$ . Membrane permeabilization was completed within 10–15 min. Under these experimental conditions, recorded currents remained unchanged for 1.5–2 h. Experiments were carried out at room temperature (20–22°C). Currents, recorded with a patch-clamp amplifier ("Biopripor," Special Construction Bureau, Pushchino, Russia) integrated with a PCL-718 LabCard and an IBM AT/386 computer, were analyzed using the BioQuest software developed by A. E. A. and Y. M. K. Average values of the recorded parameters are presented as mean  $\pm$  SE.

In cardiomyocytes from active ground squirrels, the three-dimensional current-voltage relationship, recorded as a response to a voltage-ramp at different holding potentials, possesses two distinct inward current components (see Fig. 1 A). The major component (peak 1) is sensitive to tetrodotoxin, reaches maximum at holding potentials more negative than –100 mV, and is related to the  $\text{Na}^+$  current. The minor component (peak 2) is blocked by nifedipine, and is related to the  $\text{Ca}^{2+}$  current. Because at holding potentials between –60 and –70 mV only  $\text{Ca}^{2+}$  current could be recorded (see Fig. 1 A), measurements of  $\text{Ca}^{2+}$  current in cardiocytes from ground squirrels (both in hibernating and active states) were performed in this study without the use of a blocker of  $\text{Na}^+$  channels.

### Determination of kinetic parameters

Although kinetic parameters of various ionic currents in cardiocytes are usually determined from experimental records obtained by means of special stimulation protocols (Imaizumi et al., 1989; Langton et al., 1989), in the present study kinetic parameters were determined using one-pulse rectangular stimulation. To develop a kinetic model of  $\text{Ca}^{2+}$  current we

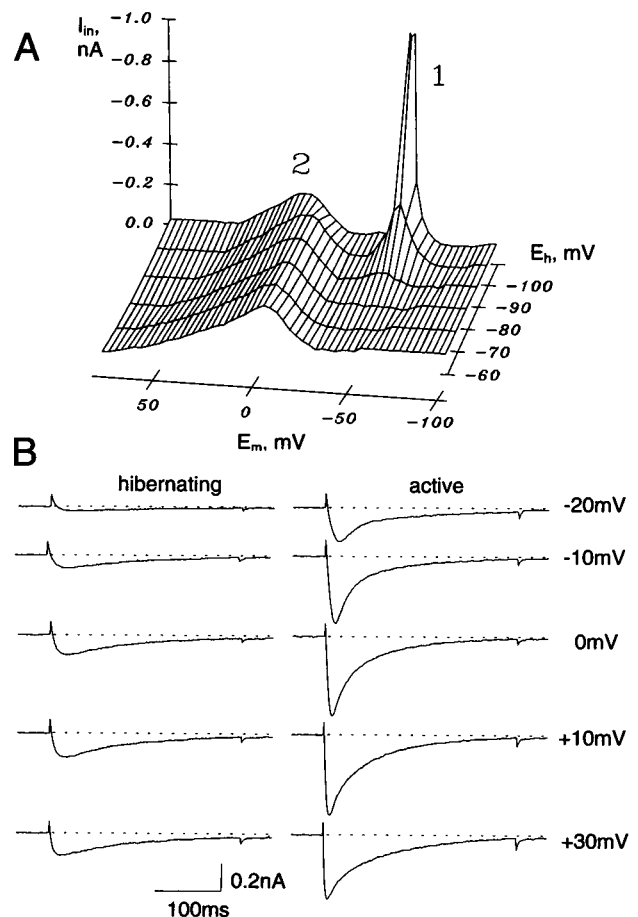


FIGURE 1 A: Three-dimensional reconstruction of the total inward current in isolated cardiocytes from active ground squirrel. Inward current ( $I_{in}$ , z axis) plotted as a function of membrane potential ( $E_m$ , x axis) and holding potential ( $E_h$ , y axis). Current-voltage relationships obtained for a voltage ramp of 1 V/s. B:  $\text{Ca}^{2+}$  currents of hibernating (left) and active (right) ground squirrels. Recordings made in response to depolarizing pulses to indicated membrane potentials (right). Holding potential, –60 mV. (.....) Zero current.

used the modified Hodgkin-Huxley equations (Hodgkin and Huxley, 1952; Adams and Gage, 1979), and approximated experimental current curves by minimization (Himmelblau, 1972) using the Nelder-Mead method of deformed polyhedron. The parameters of the  $\text{Ca}^{2+}$  current model were determined using the least square method.

## RESULTS

### Existence of a reduced $\text{Ca}^{2+}$ current in cardiomyocytes from hibernating ground squirrels

We compared voltage-dependent  $\text{Ca}^{2+}$  currents in cardiomyocytes isolated from ground squirrels during hibernation (Fig. 1 B, left) versus active state (Fig. 1 B, right) under the same experimental conditions. At each membrane potential (from –20 to +30 mV), the magnitude of the  $\text{Ca}^{2+}$  current was approximately threefold greater in active versus hibernating animals (Fig. 1 B). The average peak values of the  $\text{Ca}^{2+}$  current was  $-0.133 \pm 0.02$  nA ( $n = 26$ ) and

$-0.335 \pm 0.04$  nA ( $n = 16$ ) in hibernating and in active ground squirrels, respectively. Thus, comparative analysis revealed that the  $\text{Ca}^{2+}$  current is reduced during hibernation (see also Alekseev et al., 1994).

### Determination of kinetic parameters of $\text{Ca}^{2+}$ current in hibernators: choice of model

The kinetic parameters of the voltage-dependent  $\text{Ca}^{2+}$  current in hibernators were determined by fitting the experimentally obtained  $\text{Ca}^{2+}$  current curves with the function derived from the modified Hodgkin-Huxley equation:

$$I_{\text{Ca}} = \overline{g_{\text{Ca}}} d^r f^s (E - E_{\text{Ca}}) \quad (1)$$

$$\dot{d} = \frac{d_{\infty} - d}{\tau_d}$$

$$\dot{f} = \frac{f_{\infty} - f}{\tau_f}$$

where  $I_{\text{Ca}}$  is the  $\text{Ca}^{2+}$  current,  $\overline{g_{\text{Ca}}}$  is the maximal  $\text{Ca}^{2+}$  conductance,  $E_{\text{Ca}}$  is the equilibrium potential for the  $\text{Ca}^{2+}$  current;  $d_{\infty}$  and  $f_{\infty}$  are variables of  $\text{Ca}^{2+}$  channel activation and inactivation, and steady-state activation and inactivation parameters, which are functions of the membrane potential  $E$ ;  $\tau_d$  and  $\tau_f$  are time constants of activation and inactivation, which are also functions of  $E$ ; and  $r$  and  $s$  are the powers of the activation and inactivation variables.

To calculate the minimum  $\chi^2$ , using the Nelder-Mead method, the parameters,  $d_{\infty}$  and  $f_{\infty}$ ,  $\tau_d$ ,  $\tau_f$ ,  $r$ , and  $s$  were chosen so that the solution to the system Eq. 1 describes most closely the experimentally obtained time course of the  $\text{Ca}^{2+}$  current (Fig. 2). The solution to the differential equations of system Eq. 1 was used as a function to

calculate  $\chi^2$ :

$$I = g_{\text{Ca}}(E - E_{\text{Ca}}) \cdot \{d_{\infty} - (d_0 - d_{\infty})e^{-d_0/\tau_d}\}^r \cdot \{f_{\infty} - (f_0 - f_{\infty})e^{-f_0/\tau_f}\}^s \quad (2)$$

where  $d_0$  and  $f_0$  are the steady-state activation and inactivation current components at a given holding potential, respectively. Because the threshold of  $\text{Ca}^{2+}$  channel activation occurs at potentials more positive than the holding potential, we used the limiting values for these parameters ( $d_0 = 0$ ,  $f_0 = 1$ ).

According to the literature (e.g., Hagiwara and Nakajima, 1966),  $E_{\text{Ca}}$  in myocytes is of the order of 150 mV, and the value of  $g_{\text{Ca}}$  strongly depends on experimental conditions. Herein,  $g_{\text{Ca}}$  was determined from the slope of the right branch of the current-voltage curve for the peak current (Table 1; see also Fig. 3 A). Because experimentally obtained values of  $g_{\text{Ca}}$  were somewhat lower than the maximally achievable values because of the inactivation process, we used in most cases  $E_{\text{Ca}} = 170$  mV, which is close to the upper theoretical limit. Actually, for values around 150–160 mV, we obtained kinetic parameters that did not essentially differ from the value obtained using 170 mV as the  $E_{\text{Ca}}$  value; however, with values  $<170$  mV the fit of the experimental data with Eq. 2 (defined using the total mean-square error as the criterion) was not as adequate as with 170 mV.

The difference between the experimental and theoretical (obtained using Eq. 2) curves did not exceed a few percent (Fig. 2 A) regardless of the values for  $E$  and/or  $t$ , thus, indicating appropriate fitting. However, curves for  $d_{\infty}(E)$  and  $f_{\infty}(E)$  had a nonmonotonic shape even after varying the  $r$  and  $s$  parameters (i.e., (1, 2), (1, 4), and (2, 2)) (see Fig. 2, B–D) which makes difficult the interpretation of the activation and inactivation processes. Only by using  $r = 2$ , it was

FIGURE 2 A: Experimental (dotted) and theoretical (solid)  $\text{Ca}^{2+}$  current curves. Theoretical curve obtained by fitting experimental data by Eq. 2. (·····) Zero current. Along the zero current line, the theoretical curve is subtracted from the experimental curve. B–D: Results of fitting currents at different values of  $r$  and  $s$  by Eq. 2. (○) Steady-state solutions for activation  $d_{\infty}(E)$ ; (□) steady-state solutions for inactivation  $f_{\infty}(E)$ .

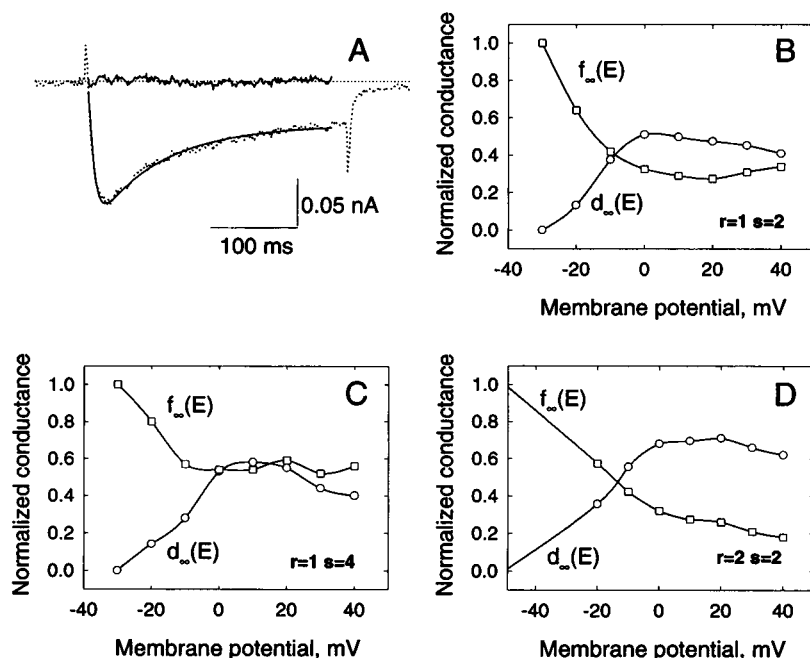


TABLE 1 Examples of kinetic parameters calculated by Eq. 4\*

$E$ , mV	$I_{\max}$ , nA	$d_{\infty}(E)$	$\tau_d$ , ms	$f_{1\infty}(E)$	$\tau_{f1}$ , ms	$f_{2\infty}(E)$	$\tau_{f2}$ , ms
-20	-0.120	0.395	5.57	0.538	96.2	0.875	0.78
-10	-0.268	0.641	6.54	0.325	76.9	0.825	31.2
0	-0.370	0.767	4.92	0.246	80.1	0.788	24.6
10	-0.393	0.790	3.64	0.174	120.2	0.686	16.9
20	-0.363	0.858	3.65	0.127	147.5	0.600	7.45
30	-0.312	0.938	3.46	0.108	159.1	0.462	4.76
40	-0.250	1.000	3.30	0.050	225.7	0.317	4.24
$g_{Ca} = 5.76 \text{ nS}$ $E_{Ca} = +170 \text{ mV}$							
$E_d^{0.5} = -16.1 \text{ mV}$ $k_d = 16.4 \text{ mV}$							
$E_{f1}^{0.5} = -20.4 \text{ mV}$ $k_{f1} = -19.7 \text{ mV}$							
$E_{f2}^{0.5} = 26.3 \text{ mV}$ $k_{f2} = -21.3 \text{ mV}$							

\*The results of fitting are exemplified in Fig. 4.

possible to approach a monotonous shape for the curve (Fig. 2 D). Thus, the kinetics of activation and inactivation of  $Ca^{2+}$  current measured in cardiocytes of hibernators cannot be represented as a monoexponential process within the frame of the model defined by the system Eq. 1. This agrees with the difficulty to describe the activation and inactivation of  $Ca^{2+}$  current as monoexponential processes (Imaizumi et al., 1989; Langton et al., 1989; Markwardt and Nilius, 1988).

To remain within the conventional interpretation of  $Ca^{2+}$  channel kinetics we modified the system Eq. 1, and involved two, instead of one, independent inactivation

processes:

$$I_{Ca} = g_{Ca} d^r f_1^{s_1} f_2^{s_2} (E - E_{Ca})$$
$$\dot{d} = \frac{d_{\infty} - d}{\tau_d} \tag{3}$$
$$\dot{f}_i = \frac{f_{i\infty} - f_i}{\tau_{fi}}, \quad i = 1, 2$$

where  $f_1$  and  $f_2$  are variables corresponding to the first and second inactivation process, respectively;  $\tau_{f1}$  and

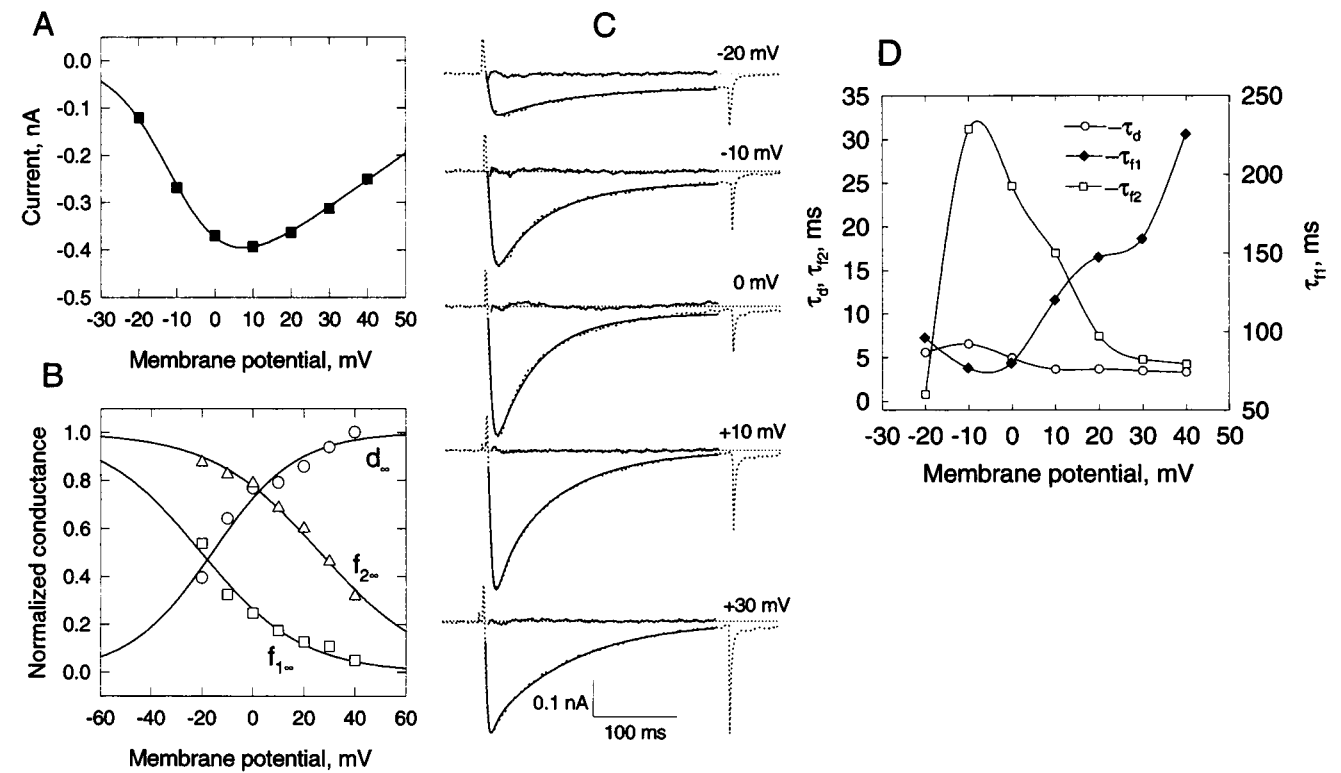


FIGURE 3 A: Current-voltage relationship constructed from the peak current values.  $g_{Ca}$ , calculated from the slope of right branch of the voltage-current relationship, is 5.76 nS. B: Normalized conductance versus membrane potential. (○) Activation; (□) slow inactivation; (△) rapid inactivation; (—) approximation by the Boltzmann equation (Eq. 5). C: Reconstruction of  $Ca^{2+}$  currents. (·····, —) Experimental and theoretical currents, respectively. Along the zero current line, the theoretical curve is subtracted from the experimental curve. D: Characteristic times of  $Ca^{2+}$  currents obtained by fitting the experimental data using Eq. 4.

$\tau_{f2}$ , the respective time constants; and  $s_1$  and  $s_2$ , the respective exponents. The solution to the system Eq. 3 was used as a function to fit the experimentally obtained curves:

$$I = g_{Ca}(E - E_{Ca}) \cdot \{d_{\infty} - (d_0 - d_{\infty})e^{-t/\tau_d}\} \cdot \prod_i \{f_{i\infty} - (f_{i0} - f_{i\infty})e^{-t/\tau_{fi}}\}^{s_i}, \quad i = 1, 2 \quad (4)$$

Using Eq. 4 to fit the experimentally obtained  $Ca^{2+}$  current curves, we calculated  $d_{\infty}(E)$ ,  $f_{1\infty}(E)$ ,  $f_{2\infty}(E)$ ,  $\tau_d$ ,  $\tau_{f1}$ , and  $\tau_{f2}$ , for currents obtained in hibernating and active ground squirrels. The best fit of the experimentally obtained values to the model (defined by the criterion of the total mean-square error for all  $E$  and  $t$  values) was obtained at  $r = 2$ ,  $s_1 = 2$ , and  $s_2 = 1$ .

Fig. 3 demonstrates that the use of Eq. 4 permits the expected dependence of all normalized conductance on the membrane potential (Fig. 3 B) and provides an insignificant discrepancy between the theoretical and experimental current curves (Fig. 3 C). Indeed, it can be seen that, through the entire range of membrane potentials, all dependencies that characterize steady-state activation and both inactivation processes possess the expected monotonic shape (Fig. 3 B). Consequently, these parameters can be well described by the Boltzmann equation:

$$\text{for activation } d_{\infty}(E) = [1 + e^{(E_d^{0.5} - E)/k_d}]^{-1} \quad (5)$$

where  $E$  is the potential at which the steady-state activation reaches its half-value, and  $k_d$  is the slope parameter. Correspondingly for inactivation,  $E_{f1}^{0.5}$ ,  $E_{f2}^{0.5}$ ,  $k_{f1}$ , and  $k_{f2}$  denote the same parameters in the Boltzmann function. Table 1 shows typical kinetic parameters obtained using Eq. 4 and fitted by the Boltzmann equation (Eq. 5).

The present model of  $Ca^{2+}$  current (developed within the frame of the Hodgkin-Huxley theory) had to incorporate an additional parameter of the inactivation process to obtain a satisfactory description of the  $Ca^{2+}$  current kinetics in hibernators. As the additional inactivation component ( $f_2$ ) had faster characteristic times when compared with the  $f_1$ -inactivation process (Fig. 3 and Table 1), we assigned  $f_2$  to denote fast, and  $f_1$  to denote slow inactivation throughout the text.

### Kinetic scheme of $Ca^{2+}$ current in hibernators

Based on the described model, the scheme of the kinetics of  $Ca^{2+}$  current in hibernators was defined. To this end, differential equations of system Eq. 3 were rewritten as:

$$\begin{aligned} \dot{d} &= \alpha_d(1 - d) - \beta_d d \\ \dot{f} &= \alpha_{fi}(1 - f_i) - \beta_{fi} f_i, \quad i = 1, 2; \end{aligned}$$

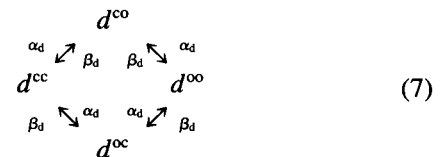
where the rate constants  $\alpha$  and  $\beta$  are related to the approx-

imation parameters as follows:

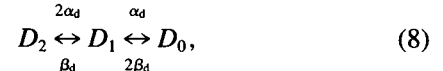
$$\begin{aligned} \alpha_d &= \frac{d_{\infty}}{\tau_d}, \quad \beta_d = \frac{1}{\tau_d} - \alpha_d; \\ \alpha_{fi} &= \frac{f_{i\infty}}{\tau_{fi}}, \quad \beta_{fi} = \frac{1}{\tau_{fi}} - \alpha_{fi}, \quad i = 1, 2 \end{aligned} \quad (6)$$

In terms of opening and closing of activation and inactivation gates, parameters  $\alpha_d$ ,  $\alpha_{f1}$ ,  $\alpha_{f2}$  and  $\beta_d$ ,  $\beta_{f1}$ ,  $\beta_{f2}$  are rate constants for the opening and closing of the activation gate and the  $f_1$ - and  $f_2$ -types of inactivation gates.

Given that, depending on the value of  $r$ ,  $s_1$ , and  $s_2$ , a channel has by convention two activation gates, two  $f_1$ -type inactivation gates, and one  $f_2$ -type inactivation gate, the scheme of activation and inactivation can be represented as:

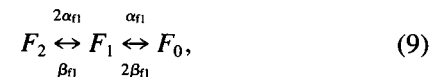


where  $d^{cc}$  is the state of the channel in which two activation gates are closed,  $d^{co}$  and  $d^{oc}$  are states in which one gate is open whereas the other is closed, and  $d^{oo}$  is the state in which both gates are open. Considering that states  $d^{co}$  and  $d^{oc}$  are macroscopically indistinguishable:



where  $D_2 = d^{cc}$ ,  $D_1 = (d^{co} + d^{oc})$ ,  $D_0 = d^{oo}$ .

Similarly, a kinetic scheme for slow inactivation ( $f_1$ -type inactivation) can be constructed. Here,  $f_1^{cc}$  denotes the state in which two gates are closed,  $f_1^{co}$  and  $f_1^{oc}$  denote the states in which one gate is closed whereas the other is open, and  $f_1^{oo}$  denotes the state in which both inactivation gates are open. Again, because states  $f_1^{co}$  and  $f_1^{oc}$  are indistinguishable, the kinetic scheme of slow inactivation can be transformed:



where  $F_2 = f_1^{cc}$ ,  $F_1 = (f_1^{co} + f_1^{oc})$ ,  $F_0 = f_1^{oo}$ .

For the  $f_2$ -type inactivation:



where  $G_1$  and  $G_0$  are the states with closed and open gates of the  $f_2$ -type, respectively. Thus, Eqs. 8, 9, and 10 compose a scheme of the  $Ca^{2+}$  current in hibernators.

Based on this scheme (Eqs. 8–10), we also constructed a more general kinetic scheme of the  $Ca^{2+}$  channel (not shown). This scheme involves 18 states of  $D_i F_j G_k$  ( $i, j = 0, 2; k = 0, 1$ ) with transitions between states occurring at rate constants defined in Eqs. 8–10 in accordance

with the principle of independence of activation and both inactivations.

To quantify the proposed model under baseline conditions, as well as after cAMP-dependent phosphorylation,  $\text{Ca}^{2+}$  currents in the absence and in the presence of a saturating concentration (Kurachi et al., 1989) of a  $\beta$ -adrenoceptor agonist, isoproterenol, were analyzed in ground squirrels, and the  $d_{\infty}(E)$ ,  $f_{1\infty}(E)$ ,  $f_{2\infty}(E)$ ,  $\tau_d$ ,  $\tau_{f1}$ ,  $\tau_{f2}$ ,  $\alpha_d$ ,  $\alpha_{f1}$ ,  $\alpha_{f2}$ ,  $\beta_d$ ,  $\beta_{f1}$ , and  $\beta_{f2}$  parameters calculated.

Kinetics of activation

The steady-state activation parameters (shown in Fig. 4 and Table 2) indicate that the normalized conductance of  $\text{Ca}^{2+}$  current (Fig. 4 A) in cardiomyocytes from hibernating and active ground squirrels are not markedly different, at least for potentials more positive than  $-10$  mV.

Addition of isoproterenol (2 mM) shifted the activation curve by 10 mV toward hyperpolarization in myocytes from hibernating animals (Fig. 4 A). A significant difference between active and hibernating animals was observed in the rate constants of activation at all examined potentials (Fig. 4 B). This indicates that stimulation of  $\beta$ -adrenoreceptors modulates the kinetics of  $\text{Ca}^{2+}$  current activation. This finding is in accord with the hypothesis that insufficient phosphorylation of  $\text{Ca}^{2+}$  channels may occur in hibernating animals. However, if the lack of phosphorylation of  $\text{Ca}^{2+}$  channels during hibernation was the sole mechanism responsible for changes in  $\text{Ca}^{2+}$

TABLE 2 Averaged steady-state characteristics of  $\text{Ca}^{2+}$  current activation and effective valences of gating charge

	Active animals	Hibernating animals	Hibernating animals + isoproterenol
$E_d^{0.5}$ , mV	$-15.6 \pm 4.7$	$-10.0 \pm 2.0$	$-19.3 \pm 3.0$
$k_d$ , mV	$17.9 \pm 1.5$	$14.6 \pm 2.4$	$16.4 \pm 1.3$
$\alpha_d^0$ , $\text{ms}^{-1}$	$0.195 \pm 0.0182$	$0.110 \pm 0.0146$	$0.204 \pm 0.0480$
$\beta_d^0$ , $\text{ms}^{-1}$	$0.088 \pm 0.0136$	$0.049 \pm 0.0073$	$0.060 \pm 0.0144$
$z_d$	$1.52 \pm 0.15$	$1.91 \pm 0.22$	$1.60 \pm 0.13$

current, then the magnitude of the shift induced by isoproterenol could be expected to be more pronounced.

The relationship between constants  $\alpha_d$  and  $\beta_d$  versus membrane potential (Fig. 4, C and D) for the activation kinetic scheme (see Eq. 8) shows that the slopes of the curves  $\alpha_d$  and  $\beta_d$  are of an opposite sign (see also Fig. 5 B). This could indicate that the process of opening and closing of activation gates may be represented as a single stage conformational transition for  $\text{Ca}^{2+}$  channels in myocytes obtained during the hibernating and active states, and can be approximated as:

$$\alpha_d = \alpha_d^0 \cdot \exp\left(z_d \Theta \frac{FE}{RT}\right)$$
$$\beta_d = \beta_d^0 \cdot \exp\left(z_d (\Theta - 1) \frac{FE}{RT}\right)$$

This equation describes conformational transition related to the movement of gating charge  $z_d$  across the membrane.

FIGURE 4 Averaged steady-state characteristics of  $\text{Ca}^{2+}$  current activation. (□) Hibernating; (○) active; (△) hibernating plus isoproterenol ( $n = 5$ ). A: Normalized conductance. Solid curves are the result of fitting data by Eq. 5 (see also Table 2). B: Characteristic times of  $\text{Ca}^{2+}$  current activation. C and D: Rate constants for kinetic scheme (Eq. 8), as calculated from Eq. 6.

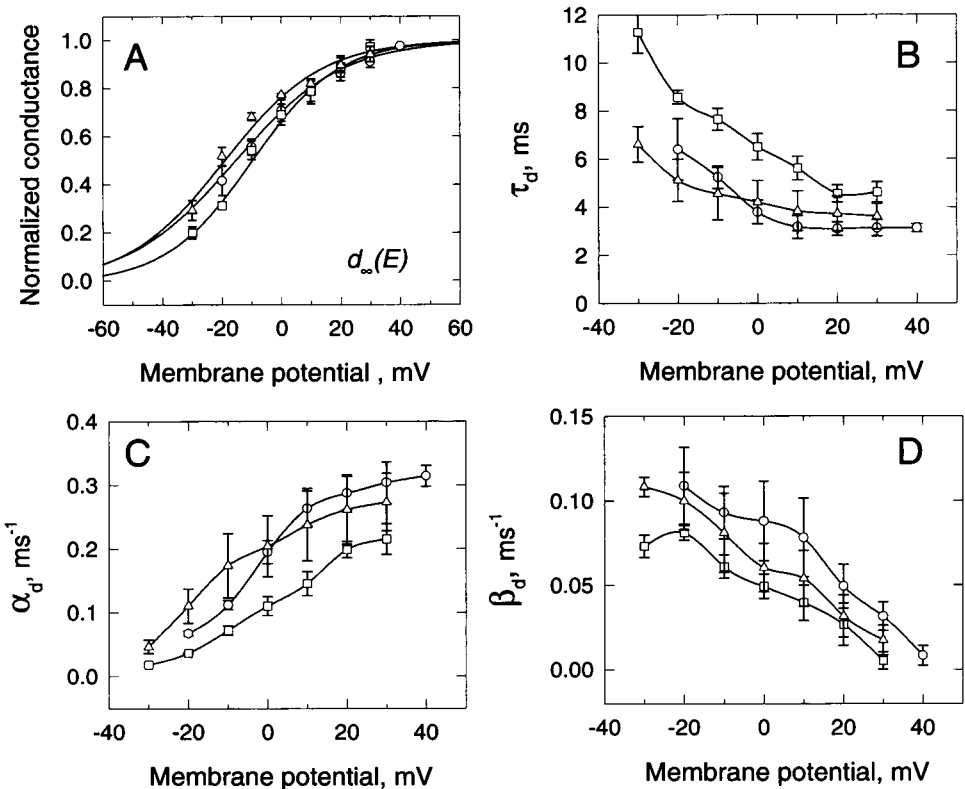
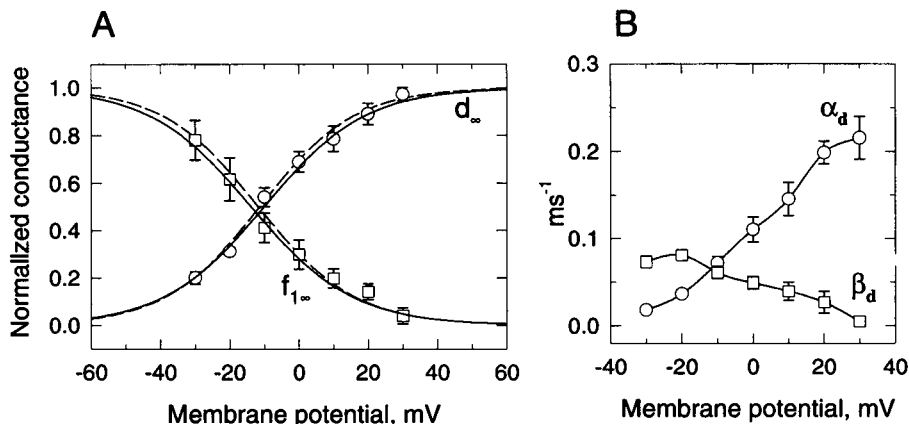


FIGURE 5 A: Normalized conductances calculated by approximating  $\text{Ca}^{2+}$  currents of hibernating ground squirrels using Eq. 4. Solid curves obtained by approximation of data using Eq. 5; dotted curves and obtained by calculation of gating charge  $z_d$  using Eq. 11. B: Averaged rate constants of kinetic scheme (Eq. 8) as calculated by fitting  $\text{Ca}^{2+}$  currents using Eq. 4.



Here  $\alpha_d^0$  and  $\beta_d^0$  are the rate constants at  $E = 0$  mV (Fig. 5 B);  $F$ ,  $R$ , and  $T$  are the constants having a general meaning;  $\Theta$  is the parameter characterizing the free energy of conformational transition, which lies between 0 and 1. The activation variable is approximated by the function:

$$d_\infty(E) = \left[ 1 + \frac{\beta_d^0}{\alpha_d^0} \exp\left(-z_d \frac{FE}{RT}\right) \right]^{-1}, \quad (11)$$

from which the effective valence of activation gating charge  $z_d$  can be calculated. The results of fitting the averaged values (Fig. 5 A) using Eq. 11 are presented in Table 2.

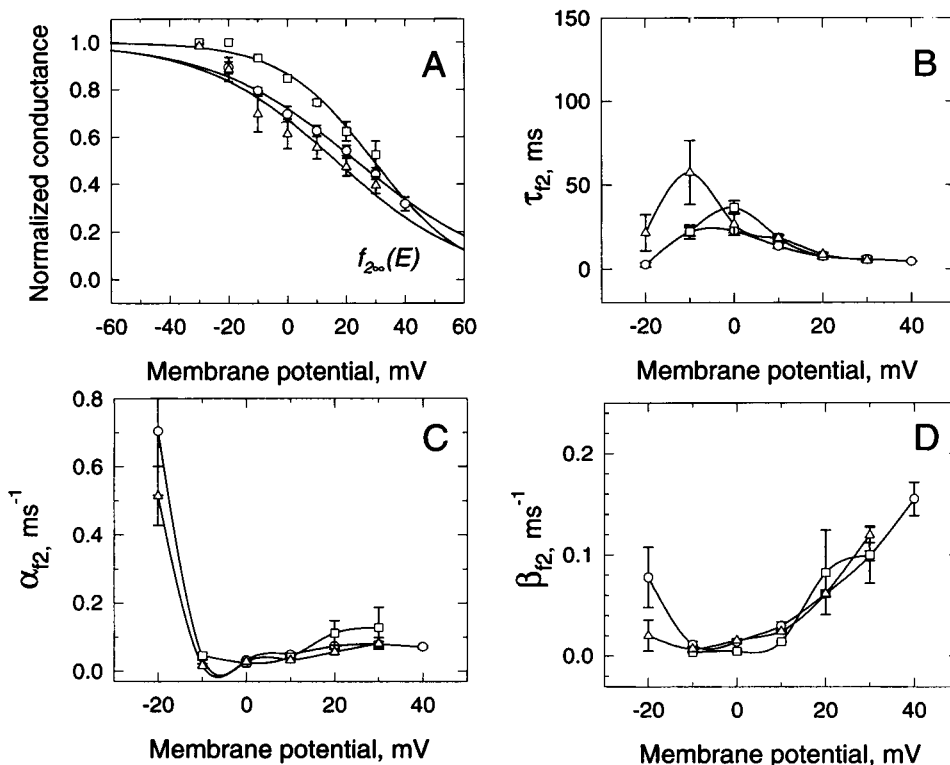
Based on the scheme of  $\text{Ca}^{2+}$  channel activation (see Eq. 9), which suggests the presence of two activation gates, the total effective charge transferred by the gate subunits on activation of  $\text{Ca}^{2+}$  channels in active ground squirrels is

estimated to be three elementary charges. This value is in good agreement with values reported for other species (Schneider and Chandler, 1973; Neely et al., 1993; Bangalore et al., 1995). Isoproterenol reversed the estimated number of elementary charges to the values calculated for myocytes from active squirrels (Table 2). Thus, at least in principle, a shift in the elementary charges, which appears to occur in hibernation, could be related to the lack of channel phosphorylation.

#### Kinetics of rapid inactivation

Fig. 6 presents averaged steady-state characteristics of rapid  $\text{Ca}^{2+}$  current inactivation in active and hibernating ground squirrels. Solid curves (Fig. 6 A) are constructed from the

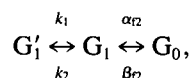
FIGURE 6 Averaged steady-state characteristics of fast  $\text{Ca}^{2+}$  current inactivation. A: Normalized conductance. Solid curves constructed using Eq. 5 (Table 3). B: Characteristic times. C and D: Constants of fast inactivation (kinetic scheme (Eq. 10)). [(squo)] Hibernating; (O) active; ( $\Delta$ ) hibernating plus isoproterenol ( $n = 5$ ).



results of approximation using the Boltzmann Eq. 5 (Table 3). A marked difference in the steepness of the inactivation curves between active and hibernating animals was obtained (Fig. 6 A). Therefore, in cardiocytes of hibernating ground squirrels, the contribution of the inactivation component to the total  $\text{Ca}^{2+}$  conductance at potentials more negative than 0–10 mV is negligible, whereas in cells of active animals it is significant (Fig. 6 A). Isoproterenol activates  $\text{Ca}^{2+}$  currents in cardiocytes of hibernating ground squirrels and shifts the normalized conductance of inactivation toward hyperpolarization and values characteristic for active animals (Fig. 6 A).

However, unlike activation, the inactivation process cannot be represented as a single-stage conformational transition related to movement of the gating charge, because the dependencies of  $\alpha_{f2}$  and  $\beta_{f2}$  (kinetic scheme, Eq. 10) on potential have similar slopes (Fig. 6, C and D).

The membrane depolarization not only shifts the equilibrium toward channel inactivation but simultaneously increases the rate of forward and reverse steps of rapid inactivation. This is indirect evidence that the inactivation process involves additional rapid potential-dependent steps. Thus, an anomalous decrease in the rate of exit from rapid inactivation on depolarization, i.e., an increase in  $\alpha_{f2}(E)$ , suggests that for rapid inactivation, not scheme Eq. 10 but the following scheme is valid:



where  $G'_1$  is the state with closed inactivation gates, which cannot change to state  $G_0$ . Provided that  $k_1$  and  $k_2$  are much greater than  $\alpha_{f2}$  and  $\beta_{f2}$ , i.e., states  $G'_1$  and  $G_1$  are related by the equilibrium relation  $G'_1 = k_1 \cdot G_1/k_2$ , the equation for rapid inactivation of system Eq. 3 is then transformed to:

$$\dot{f}_2 = \frac{\alpha_{f2}}{1 + k_1/k_2} (1 - f_2) - \beta_{f2} f_2.$$

Hence, it follows that if the ratio  $k_1/k_2$  sharply decreases on membrane depolarization, the effective rate constant of exit from inactivation  $\alpha_{f2}/(1 + k_1/k_2)$  may increase.

There is another explanation for the unusual dependence of inactivation rate constants on potential, which is not related to the movement of inactivation gating charge. Studies of gating behavior of  $\text{Na}^+$  (Gonoi and Hille, 1987; Cota and Armstrong, 1989),  $\text{K}^+$  (Koren et al., 1990) and T-type  $\text{Ca}^{2+}$  channels (Chen and Hess, 1990) showed that transi-

tion of the channel to the inactivation state is virtually independent of membrane potential; i.e., the inactivation gating charge, if any, is very small. It may be assumed that in our case the dependence of inactivation on potential is a consequence of potential-dependent channel activation.

On the basis of the analysis of the effect of nifedipine on  $\text{Ca}^{2+}$  current, we propose a mechanism for  $\text{Ca}^{2+}$  channel inactivation (Fig. 7). This antagonist of L-type  $\text{Ca}^{2+}$  channels affects only the kinetics of activation and slow inactivation (Fig. 7, top and middle), leaving rapid inactivation unaltered (Fig. 7, bottom). The characteristic time of rapid inactivation  $\tau_{f2}$  is insensitive to nifedipine. Thus, the effect of isoproterenol on activation and the anomalous dependence of  $\alpha_{f2}$  and  $\beta_{f2}$  suggests that kinetics of rapid inactivation are mainly due to cAMP-dependent phosphorylation. It is unlikely that rapid inactivation is related to the activation process, because nifedipine-induced changes in activation kinetics should also have affected the kinetics of activation. It is also questionable that the inactivation proceeds by a current-dependent mechanism, given that  $\text{Ca}^{2+}$

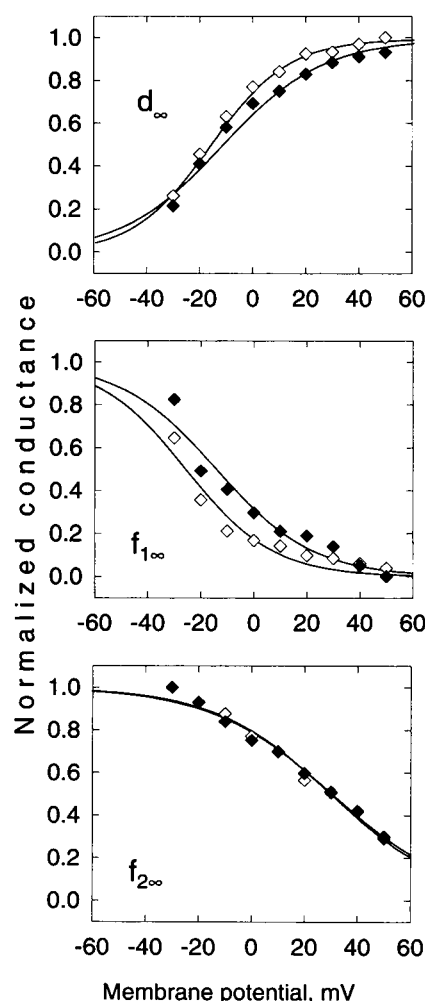


FIGURE 7 Normalized conductances calculated by fitting  $\text{Ca}^{2+}$  currents in nifedipine-treated cardiocyte using Eq. 5.  $\diamond$ , without 5  $\mu\text{M}$  nifedipine;  $\blacklozenge$ , with 5  $\mu\text{M}$  nifedipine.

TABLE 3 Averaged steady-state characteristics of  $\text{Ca}^{2+}$  current inactivation

	Active animals	Hibernating animals	Hibernating animals + isoproterenol
$E_{f1}^{0.5}$ , mV	$-26.3 \pm 3.9$	$-13.6 \pm 4.9$	$-39.7 \pm 3.9$
$k_{f1}$ , mV	$-24.0 \pm 2.6$	$-14.7 \pm 1.1$	$-24.2 \pm 2.1$
$E_{f2}^{0.5}$ , mV	$23.1 \pm 1.6$	$29.1 \pm 2.9$	$16.4 \pm 5.0$
$k_{f2}$ , mV	$-24.6 \pm 2.0$	$-15.7 \pm 1.3$	$-22.5 \pm 0.8$



conductance, as determined by the slope of the right-hand branch of the current-voltage curve, was 2.87 nS in control and 0.99 nS in the presence of nifedipine. The peak current values also differ by threefold, and hence, the amount of  $\text{Ca}^{2+}$  entering across the channels decreases in proportion. However, as mentioned above, the normalized conductance of rapid inactivation remains unaffected. Presumably, the mechanism of current-dependent inactivation is realized in the case of slow inactivation, which will be discussed in the next section.

### Kinetics of slow inactivation

The steady-state characteristics of slow inactivation kinetics and the results of Boltzmann approximation of normalized conductance are presented in Fig. 8 and Table 3. Slow inactivation is particular to late stages of  $\text{Ca}^{2+}$  currents, and characterized by long characteristic times and low rates of opening and closing of gates ( $\alpha_{f1}$  and  $\beta_{f1}$  from kinetic scheme (Eq. 9)).

From Fig. 8 it can be seen that parameters of slow  $\text{Ca}^{2+}$  current inactivation in cardiocytes of active ground-squirrels, are shifted toward hyperpolarization by  $\sim 10$  mV relative to those of hibernating animals. Addition of isoproterenol to cells of hibernating animals caused a significant (by  $\sim 20$  mV) shift of the  $E_d^{0.5}$  toward hyperpolarization.

The shape of the dependencies  $\alpha_{f1}$  and  $\beta_{f1}$  on membrane potential varies with the state of animals (Fig. 8, C and D). In cardiocytes of active and hibernating animals in the presence of isoproterenol, both parameters decrease with increasing membrane potential; the curves have a similar

slope, indicating that the process of slow inactivation cannot be represented as a simple single-stage conformational transition associated with movement of the gating charge. By contrast, in cardiocytes of hibernating animals, parameters  $\alpha_{f1}$  and  $\beta_{f1}$  show opposite directed slopes.

Fig. 9 presents a typical experiment ( $n = 4$ ) on the effect of isoproterenol on cardiocytes of hibernating animals. Isoproterenol increased the rate of the reverse transition  $\beta_{f1}$  approximately two- to threefold and changed the direction of the slope depending on membrane potentials. In terms of the conception of independence of activation and inactivation proposed in this study, this finding may be interpreted to indicate that slow inactivation represents both a potential- and current-dependent process. If the inward current is considerably (severalfold) reduced during hibernation or in response to nifedipine, the current-dependent inactivation does not affect the general inactivation process. In this case, only the first, the potential-dependent, inactivation component is retained. Interestingly, the inactivation gating charge calculated by fitting the data using Eq. 11 (Fig. 5 A) was much higher than it might be expected ( $z_{f1} = -1.83 \pm 0.2$ ;  $n = 3$  cardiocytes from hibernating animals). As in the case of activation, the effective gating charge is two times as great, since the model suggests the availability of two inactivation gates. An indirect confirmation of this interpretation is the fact that any agonist of the  $\text{Ca}^{2+}$  current shifts, as a rule, the steady state of slow inactivation toward hyperpolarization, because of initiation of a current-dependent process, whereas all antagonist (e.g., nifedipine, Fig. 7, middle) cause opposite directed changes in this parameter. Moreover, the rate constant of slow inactivation introduced

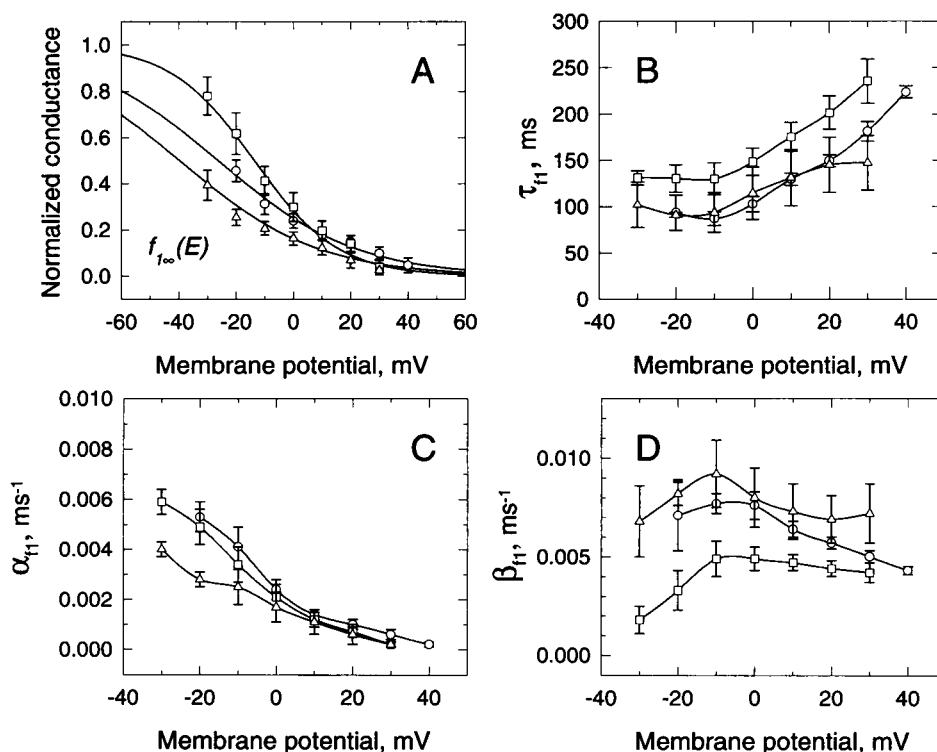
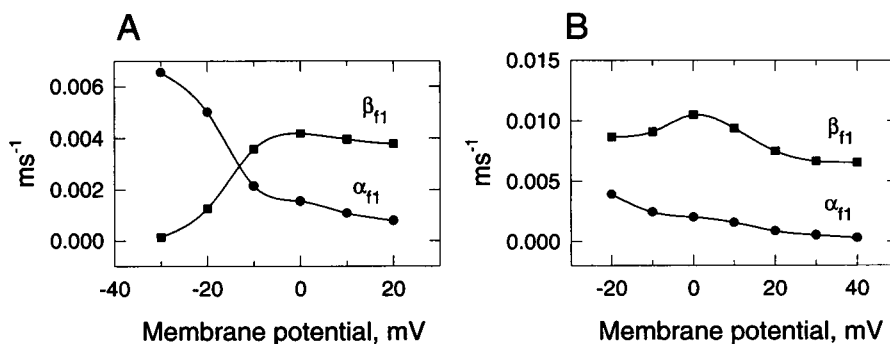


FIGURE 8 Averaged normalized conductance of slow  $\text{Ca}^{2+}$  current inactivation. A: Normalized conductance. Solid curves constructed using Eq. 5 (Table 3). B: Characteristic times. C and D: Rate constants of slow inactivation (kinetic scheme (Eq. 9)). ( $\square$ ) Hibernating; ( $\circ$ ) active; ( $\triangle$ ) hibernating plus isoproterenol ( $n = 5$ ).

FIGURE 9 Rate constants of slow  $\text{Ca}^{2+}$  current inactivation (kinetic scheme (Eq. 9)) in cardiocytes of a hibernating ground squirrel in the absence (A) and presence of isoproterenol (B).



in the model of current-dependent inactivation (Standen and Stanfield, 1982; Markwardt and Nilius, 1988) is  $0.005 \text{ ms}^{-1}$ , which coincides almost completely with the slow inactivation rate constant in our model. To verify this hypothesis, experiments involving the determination of slow inactivation parameters at different extracellular  $\text{Ca}^{2+}$  concentrations are needed. Such experiments were not carried out in the framework of this study.

## DISCUSSION

The present study investigated the properties of the L-type  $\text{Ca}^{2+}$  current in hibernators and developed a model that includes two activation gates, two  $f_1$ -type inactivation gates, and one  $f_2$ -type inactivation gate ( $d^2f_1f_2$  model). The proposed kinetic model of the L-type calcium channel with its 18 stages appears rather extensive. This is primarily because of the principle of independence of the activation and inactivation processes on which this model is based. We could have, for instance, abandoned the principle of independence of the activation and inactivation processes, as was done by Chen and Hess (1990) in the description of the gating mechanism of the T-type calcium currents. This would have reduced the number of stages in the kinetic model, but would have made the mathematical description of  $\text{Ca}^{2+}$  current kinetics much more difficult and the interpretation of experimental data less obvious. It is for this reason that we considered dissecting the process of channel regulation into relatively independent stages.

The model used for analysis of the kinetic parameters appears to be adequate. In particular, the quadratic variable of activation,  $d^2$ , indicates the existence of two closed states, which correspond to the C1-C2-0 scheme of  $\text{Ca}^{2+}$  channel activation confirmed by experiments on single  $\text{Ca}^{2+}$  channels (Tsien et al., 1986). The effective valence of the charge related to the opening of one activation gate is equal to 1.5–2, which is in good agreement with measurements of gating currents related to the activation of cardiac L-type calcium channel (Neely et al., 1993; Bangalore et al., 1995). The entire activation process is associated with movement of the total gating charge (3–4) and is sensitive to a  $\beta$ -adrenoceptor agonist. Interesting, in our opinion, is a variation in the effective valence of activation gates from

1.5 (active ground squirrels) to 2 (hibernating animals). A possible reason for this variation is the loss of one or several negatively charged phosphate groups because of dephosphorylation, assuming that the phosphorylation site is in some way connected with the gate subunits (Tsien et al., 1988). It is also of interest that  $z_d$  of  $\text{Ca}^{2+}$  channels of hibernating ground squirrels returns to values close to 1.5 in response to isoproterenol (Table 2). Presumably, the modification of this conductance component is due only to a decrease in the number of phosphorylated channels induced by a reduction in catecholamine level during hibernation (Gainullin and Saxon, 1988; Florant et al., 1982).

The slow inactivation ( $f_1$ -type) is due to two processes, one of which is related to the gate mechanism and the other to the current-dependent inactivation. For hibernating animals, the total effective gating charge lies in the range between  $-3.5$  and  $-4$ . In studies of  $\text{Ca}^{2+}$  channel kinetics, it is common to employ  $\text{Ba}^{2+}$ . It can be assumed that in our case, the function of rapid inactivation,  $f_2$ , in Eq. 3 is due to  $\text{Ca}^{2+}$ -induced inactivation of the  $\text{Ca}^{2+}$  channel. When we replaced  $\text{Ca}^{2+}$  by 2 mM  $\text{Ba}^{2+}$ , currents across  $\text{Ca}^{2+}$  channels showed a retarded inactivation. However, in this case we also failed to derive monotonic steady-state characteristics of conductance in the framework of Eq. 1, which suggests the existence of one inactivation component. Similarly, Markwardt and Nilius (1988) showed that although  $\text{Ba}^{2+}$  has a lower affinity for the calcium-binding site of the channels than  $\text{Ca}^{2+}$ ,  $\text{Ba}^{2+}$  is involved in the inactivation of  $\text{Ca}^{2+}$  channels. It is by this mechanism the two-stage course of inactivation was explained. The rate constants for rapid and slow inactivation used in their calculations were  $0.8 \text{ ms}^{-1}$  and  $0.005 \text{ ms}^{-1}$ , respectively. The rate constants derived with our model are consistent with these values (Figs. 6, 8, 9). Hadley and Lederer (1991) showed that Ca-dependent inactivation is sensitive to  $\beta$ -adrenergic stimulation. In terms of the proposed model, we were able to separate the inactivation component. This conductance component was found to be completely insensitive even to a threefold decrease in inward  $\text{Ca}^{2+}$  current induced by nifedipine.

Differences in the current kinetics between active and hibernating animals were observed for all components. In hibernating ground squirrels, the most pronounced shift toward depolarizing potentials was observed for the steady-

state conductance of both inactivation components, whereas on activation the shift was smaller. The characteristic times ( $\tau_d$ ,  $\tau_{f1}$ ,  $\tau_{f2}$ ) for all three components of  $\text{Ca}^{2+}$  current conductance in hibernating ground squirrels were found to be 1.5 times smaller than in active animals.

Isoproterenol reversed the characteristic times to values close to those of active animals, shifting simultaneously the normalized conductance toward polarizing potentials by values exceeding the difference in conductance between active and hibernating ground squirrels. This observation appears to be paradoxical to some extent, because the restoration of the kinetic parameters induced by phosphorylation of channels in hibernating animals does not govern the increase in  $\text{Ca}^{2+}$  current amplitude to values recorded in active animals in response to isoproterenol (Alekseev et al., 1994). Therefore, the change in the amplitude of  $\text{Ca}^{2+}$  current during hibernation may be due to a seasonal decrease in the single channel conductance and/or in the number of  $\text{Ca}^{2+}$  channels activated on depolarization during hibernation. Another possibility of suppression of the  $\text{Ca}^{2+}$  current may be related to the cGMP-dependent channel phosphorylation (Wahler et al., 1990; Tohse and Sperelakis, 1991). However, the role of the cGMP-dependent protein kinase in hibernating cycles requires additional investigation. Thus, the question related to mechanisms responsible for the decrease in the number of  $\text{Ca}^{2+}$  channels and/or the maintenance of these channels in a non-active state remains open.

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